

Stereocontrolled photoinduced electron transfer in metal-fullerene hybrids.

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Abstract

Photoinduced electron transfer in transition-metal complexes linked to a fullerene moiety is of increasing interest. Recently, several stereoisomers of an Ir-complex exhibiting configurational stability at metal center, which does not undergo epimerization have been synthesized (*Angew. Chem.* 56, 2136). The presence of multiple electron donor and acceptor sites located at opposite ends with respect to metallic center, creates the prerequisites for the formation of entirely different CT states. Here we report the results of quantum mechanical calculation and detailed analysis of excited-state properties for all stereoisomers of the junction. We found that the stereoisomers demonstrate clearly different charge transfer (CT) properties by photoexcitation. The found photo-stereospecific effects can be used to design new hybrids with different type of photoinduced CT states exhibiting dissimilar activity in photocatalysis.

Introduction

Conversion of the sunlight into chemical potential is considered to be one of the most important and fundamental chemical reactions in nature.^{1,2} It is based on inter- or intra-molecular electron transfer between donor and acceptor parts of the system caused by photoexcitation. Transition-metal complexes of Ru(II),^{3,4} Re(I),^{5,6} and Ir(III)^{6,7} are among most interesting systems exhibiting photoinduced electron transfer (PET). They are used for example in lighting devices,^{8,9} solar cells,¹⁰ photoredox catalysis,^{11,12} and sensors.^{13,14} The unique chemical and physical properties of fullerene C₆₀,¹⁵⁻¹⁸ such as its highly efficient electron acceptor character and electronic absorption throughout the UV-Vis spectral region, have aroused considerable interest in its application in intramolecular PET reactions¹⁹⁻²¹ and make hybrid systems containing both transition-metal complex and fullerene moiety to be of great interest.²²⁻²⁴ In particular, chiral metal complexes provide a promising tool in organometallic chemistry due to their pertinence to enantioselective reactions and catalytic applications in organic synthesis.^{25,26} Recent discovery of fullerene C₆₀ half-sandwich transition-metal hybrid complexes with chiral centers both on ligand and metal species opens new opportunities not only for metal catalysis but also for PET reactions.²⁴ The fact that metallic center in Ir pyrrolidino[3,4:1,2][60]fullerene half-sandwich complexes is not subjected to epimerization provide ample opportunities for using these complexes in photocatalytic

reactions and therefore the investigation of the nature of excited states for its stereoisomers is essential for understanding PET reaction.

Herein, we report an exhaustive DFT computational study of Ir pyrrolidino[3,4:1,2][60]fullerene half-sandwich complexes with four stereocenters: Ir-metal center, the asymmetric nitrogen atom and two C5 and C2 chiral carbon atoms at the pyrrolidine ring (Figure 1). The stereoisomers demonstrate different photoinduced electron transfer properties.

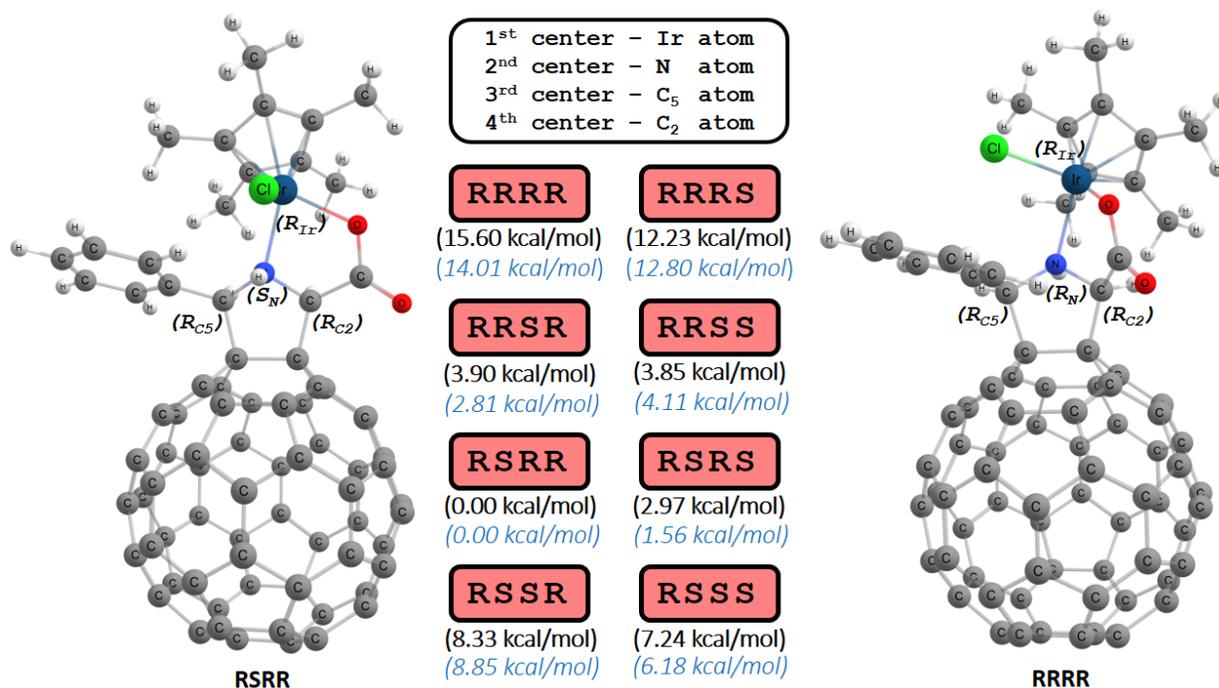


Figure 1. The most stable (RSRR) and least stable (RRRR) stereoisomers of studied Ir-complex as well as relative energies in lowest singlet (bold) and triplet (italic) states for studied stereoisomers, obtained at BLYP-D3(BJ)/def2-SVP level of theory.

Results

Stability of the stereoisomers. For complex containing four optical centers, there are sixteen possible stereoisomers or eight pairs of enantiomers. Because physical properties of enantiomers are indistinguishable (except for the interaction with polarized light), we consider only 8 stereoisomers. DFT calculations at BLYP-D3(BJ)/def2-SVP level of theory were carried out to gain more insight into the relative stability of the considered isomers in the singlet ground state (S_0). Calculations revealed that all isomers belong to the less than 16 kcal/mol energetic range (Figure 1). Studied set of structures can be divided into two groups: low-lying (RSRR, RSRS, RRSS, and RRSR) and high-lying (RSSS, RSSR, RRRS, and RRRR) isomers. In the lowest triplet state (T_1), energy profile for the studied isomers is very similar to the singlet one. RSRR isomer exhibit the greatest stability, while RRRR is less stable with the energy difference between them of about 14 kcal/mol. Regardless of isomer, the singlet state is the lowest-lying state. The S_0 - T_1 gap is characterized by a value of 28-30 kcal/mol depending on the isomer (detailed data about relative stabilities of isomers in S_0 and T_1 states are given in Table S1 and Figure S1, SI). According to Martin and co-workers,²⁴ relative stability of isomers is determined by two factors acting in different directions:

(1) strong repulsive interactions between lone pair of Cl ligand and π -system of phenyl ring and (2) stabilizing CH- π interactions between CH₃ group of methylated cyclopentadienyl ligand and π -system of phenyl ring.²⁷

The studied complexes are unique objects. From one side, four chiral centers open a variety of potential opportunities in the field of (stereo)chemical transformations.²⁶ From the other side, a chemical structure of complexes itself, namely the fact that electron acceptors (fullerene C₆₀ and Cp*) are located at opposite ends with respect to metallic center, creates the prerequisites for the formation of entirely different CT states with localization of the exciton on different and remote from each other fragments.

Usually for analyzing photoinduced electron transfer (ET) one uses two fragment model consisting of donor and acceptor sides. In this study, a multi-fragment model is applied to examine in detail the role of metal atom and different ligands in ET. Each system has been split into 6 fragments – Ir and Cl atoms, pyrrolidine moiety with the carboxyl group, as well as phenyl (Ph), pentamethylcyclopentadienyl (Cp*) and fullerene C₆₀ fragment (Figure 2).

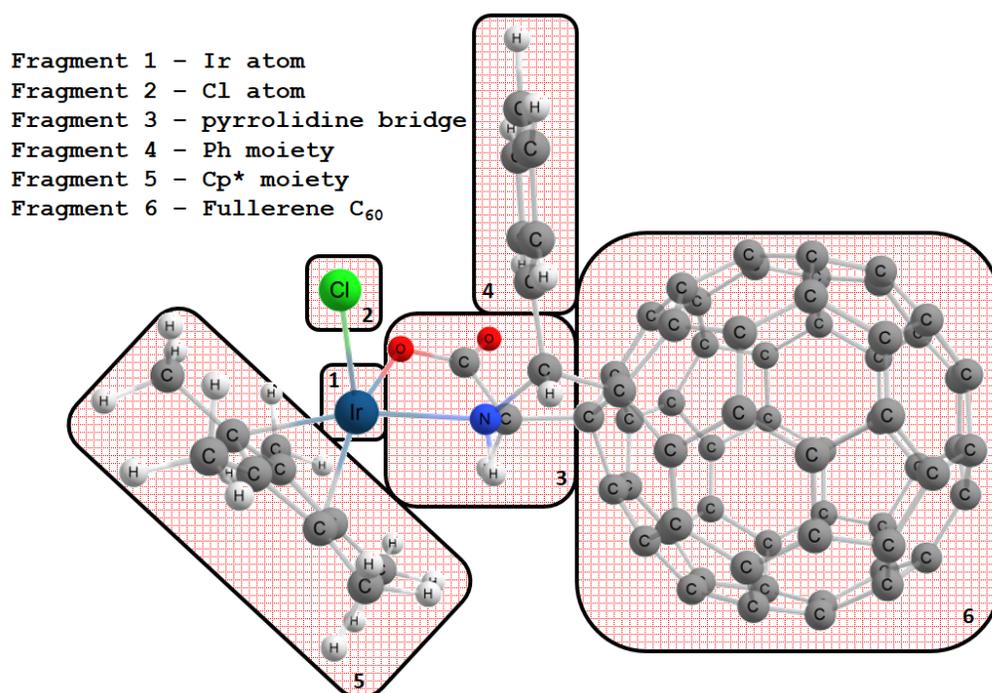


Figure 2. Fragmentation scheme demonstrated by example of RRRR isomer.

Cp* and fullerene fragments exhibit electron acceptor properties, whereas Ir, Cl, Ph and pyrrolidine bridge fragments are electron donors of different power. Analysis of excited states was carried out in terms of excitation delocalization, charge transfer (CT) and charge separation (CS) contributions. In the studied complexes, several types of excited states can be distinguished: locally excited states (LE), where exciton is mostly localized on a single fragment (in our systems all low-lying LE states are confined on C₆₀), excited states corresponding to CT and mixed states, having comparable contributions of LE and CT. Below we considered 50 lowest singlets and triplets excited states for each of isomers.

Singlet states. Predicted 50 lowest singlet excited states for each of eight studied complexes belong to the energy band from 2.4 to 4.6 eV (Table 1) in the gas-phase. Several lowest excited states are LE states

with electronic transitions occurring on C₆₀ fragment. For all isomers, first excited state is mostly associated with HOMO-LUMO transitions. First states demonstrating charge separating characteristics have been found in 3.0 – 3.2 eV region (Table 1) dependent on isomer. Specified CT states are not pure with respect to the involved fragments, *i.e.* two fragments simultaneously acts as an electron density donor. In each case, denoted states correspond to the transition of the electron density from metallic center including Cl ligand to Cp* moiety. Further, in the energy range of 4.2-4.6 eV the CT states where electronic transitions metallic center including Cl ligand to fullerene C₆₀. In the range of 4.4 – 4.6 e, electronic transitions from fragment 3 and 4 to fullerene C₆₀ are occurred.

Table 1. Singlet excitation energies (E_x, eV) major orbital contributions (HOMO(H) – LUMO(L)) and its weight, oscillator strength (f) and charge transfer values (CT, e), charge values on Ir atom in ground state (q(Ir) GS) and changes in Ir charge state (Δq Ir, e) for the studied stereoisomers of Ir half-sandwich fullerene C₆₀ complex in the gas-phase. For each transition involved fragments or fragments whose contribution is maximal (Assoc. fragments) are shown.

	Ir complex							
	RRRR	RRRS	RRSR	RRSS	RSRR	RSRS	RSSR	RSSS
	LE ₁							
E _x	2.396	2.456	2.398	2.457	2.462	2.416	2.457	2.411
Transition (weight)	H – L (0.82)	H – L (0.90)	H – L (0.85)	H – L (0.90)	H – L (0.89)	H – L (0.85)	H – L (0.90)	H – L (0.86)
f	0.006	0.004	0.006	0.004	0.005	0.006	0.004	0.006
CT	0.036	0.020	0.034	0.019	0.020	0.035	0.022	0.031
Assoc. fragments	F6	F6	F6	F6	F6	F6	F6	F6
	CT _(F1→F5)							
E _x	3.145	3.209	3.220	3.114	3.219	3.099	3.100	3.034
Transition (weight)	H-5-L+7 (0.25)	H-4-L+8 (0.34)	H-4-L+8 (0.31)	H-4-L+7 (0.31)	H-4-L+7 (0.68)	H-4-L+7 (0.26)	H-4-L+7 (0.51)	H-4-L+7 (0.44)
f	0.012	0.004	0.008	0.003	0.004	0.013	0.004	0.007
CT	0.629	0.630	0.629	0.530	0.632	0.302	0.613	0.632
Assoc. fragments	F1→F5	F1→F5	F1→F5	F1→F5	F1→F5	F1→F5	F1→F5	F1→F5
	CT _(F1→F6)							
E _x	4.376	4.303	4.605	4.214	4.345	4.297	4.256	4.449
Transition (weight)	H-5-L+1 (0.15)	H-14-L+1 (0.18)	H-6-L (0.34)	H-4-L (0.38)	H-4-L (0.39)	H-4-L (0.28)	H-4-L+1 (0.13)	H-L+5 (0.10)
f	0.031	0.019	0.040	0.015	0.013	0.015	0.002	0.050
CT	0.401	0.125	0.488	0.496	0.745	0.555	0.254	0.236
Assoc. fragments	F1→F6	F1→F6	F1→F6	F1→F6	F1→F6	F1→F6	F1→F6	F1→F6
	Most bright CT							
E _x	4.479	4.482	4.427	4.498	4.425	4.436	4.436	4.613
Transition (weight)	H-1-L+3 (0.18)	H-1-L+3 (0.14)	H-1-L+4 (0.15)	H-1-L+3 (0.18)	H-4-L+1 (0.10)	H-10-L+2 (0.12)	H-1-L+3 (0.09)	H-2-L+5 (0.26)
f	0.167	0.110	0.143	0.086	0.102	0.101	0.122	0.087

CT	0.174	0.322	0.183	0.105	0.315	0.175	0.314	0.171
q(Ir) GS	0.389	0.393	0.394	0.402	0.393	0.380	0.376	0.382
Δq Ir	0.036	0.069	0.043	0.010	0.111	0.048	0.066	0.031
Assoc. fragments	F4→F6	F3→F6	F3→F6	F3→F6	F1→F6	F3→F6	F3→F6	F3→F6

Thus, for the described system, depending on the excitation energy it is possible to populate charge transfer states with the different character. This can have a significant impact on the behavior of the complex in photoinduced catalytic reactions. However, transitions between molecular orbitals corresponding to the aforementioned (F1 → F5 and F1 → F6) CT states are characterized by the relatively small probability for the mentioned transitions. On average, oscillator strength for F1 → F5 transitions varies from 0.001 to 0.01, while for F1 → F6 it is from 0.01 to 0.03. Low oscillator strength values limit potential usage of this complexes. Then, we decided to check out the “brightest” transitions associated with CT states. For this complex, the strongest oscillator strength is about 0.09 – 0.17, dependently on particular isomer. Surprisingly, we have found that only for RSRR isomer most intense transition is relate to CT state in which electron density is transferred from Ir to fullerene C₆₀. In every other cases, transitions with highest oscillator strength were connected with pyrrolidine bridge (F3) or phenyl ring (F4) to fullerene moiety (F6) (Table 1). Visualization of the described above locally excited states and CT states characterized by F1 → F5 as well as F1 → F6 transitions for RSRR isomer are shown in Figure 3.

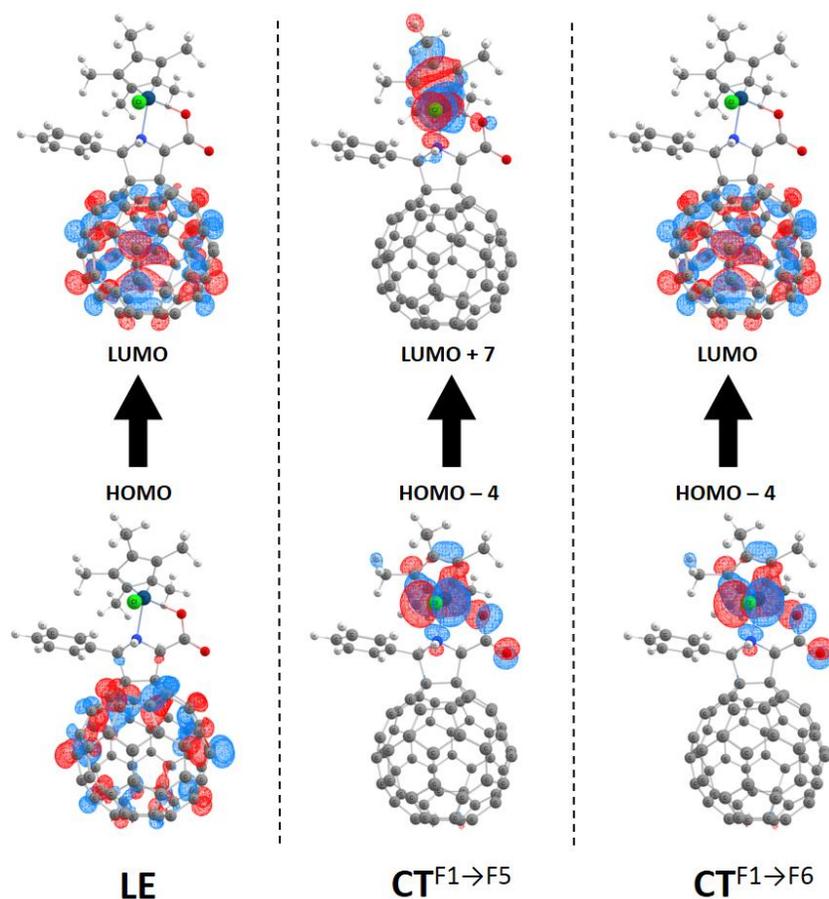


Figure 3. Frontier molecular orbitals representing lowest-lying single LE state (left), CT states correspond to F1 → F5 (middle) and F1 → F6 (right) transitions for RSRR isomer.

The electron density changes observed on Ir atom clearly demonstrate the peculiarity of RSRR stereoisomer. In accordance with Mulliken population analysis, the change of Ir charge in case of RSRR isomer is 0.11 e, which is almost twice larger than for any other isomer. At the same time, taking into account that charge on Ir atom in the ground state is about 0.38-0.40 e, the indicated changes can be considered as significant.

In other words, depending on the isomers we can efficiently generate CT excited states of different nature. Such stereospecificity towards photoinduced electron transfer provides an attractive pathway for photocatalytic reactions, where only one of the stereoisomers demonstrate activity while others remain inactive.

Triplet states. First 50 triplet excited states have been studied as well. Low-lying triplet LE^T_1 and CT^T states show behavior similar to those in the corresponding singlet states. Namely, first triplet excited state is mostly associated with the excitation localized at fullerene C_{60} moiety and corresponds to HOMO to LUMO+1 transition, while HOMO-LUMO excitation leads to the 3rd triplet LE. Similarly to singlet excited states the lowest charge separated state corresponds to the transition of the electron density from metallic center including Cl ligand to Cp^* moiety. In contrast to singlet excited states, no one of considered triplet CT excited states does not correspond to electronic transitions from Ir fragment to C_{60} . Detailed information concerning triplet excited states as well as frontier molecular orbitals representing LE^T and CT^T states are given in Table S2 and Figure S2, SI.

Solvent effect. It is well known that solvation may significantly influence both ground and excited states. Usually, the effect of solvation is relatively weak for LE states, while charge transfer states can strongly be stabilized. To assess the solvent effects on the excitation energies the equilibrium solvation model with dichloromethane (DCM) as a solvent has been applied. In the ground state, all of eight studied isomers demonstrate comparable dipole moment in the range from 7.6 to 9.6 D, depending on particular isomer. However, despite the fact that dipole moment can vary up to 20% the changes in solvation energies are only of a few percent (Table S3, S4, SI). The reason for this behavior is apparently the high steric hindrance of the complex due to large ligands that make it difficult for a solvent to access significant parts of the complex. In case of the excited states, the overall picture is very similar. Additionally, taking into account that excited states of interest is mostly associated with metallic center – the most shielded fragment of the complex – solvation energies relatively insensitive to the changes in charge state on this fragment. Comparison between excitation energies computed for the gas phase and solution indicates that LE transition energies remain almost unchanged, in spite of solvation shift for each GS and LE states is noticeable (about 0.7 eV). Detailed data for all studied complex are provided in Table S3, S4 and Figure S3, SI. Due to the already mentioned reason, CT states demonstrate untypically small shift which is comparable to that for LE states (about 0.7–0.8 eV, in case of RSRR isomer the shift is slightly more than 1.0 eV), despite the fact that significant changes in dipole moments (up to 25 D) for CT states have been observed. The effect of solvation on triplet excited states, in general, is pretty similar to those observed in singlet excited states. Figure 4 demonstrates LE and CT energies for singlet and triplet excited states both in vacuum and CPCM (DCM) solvent.

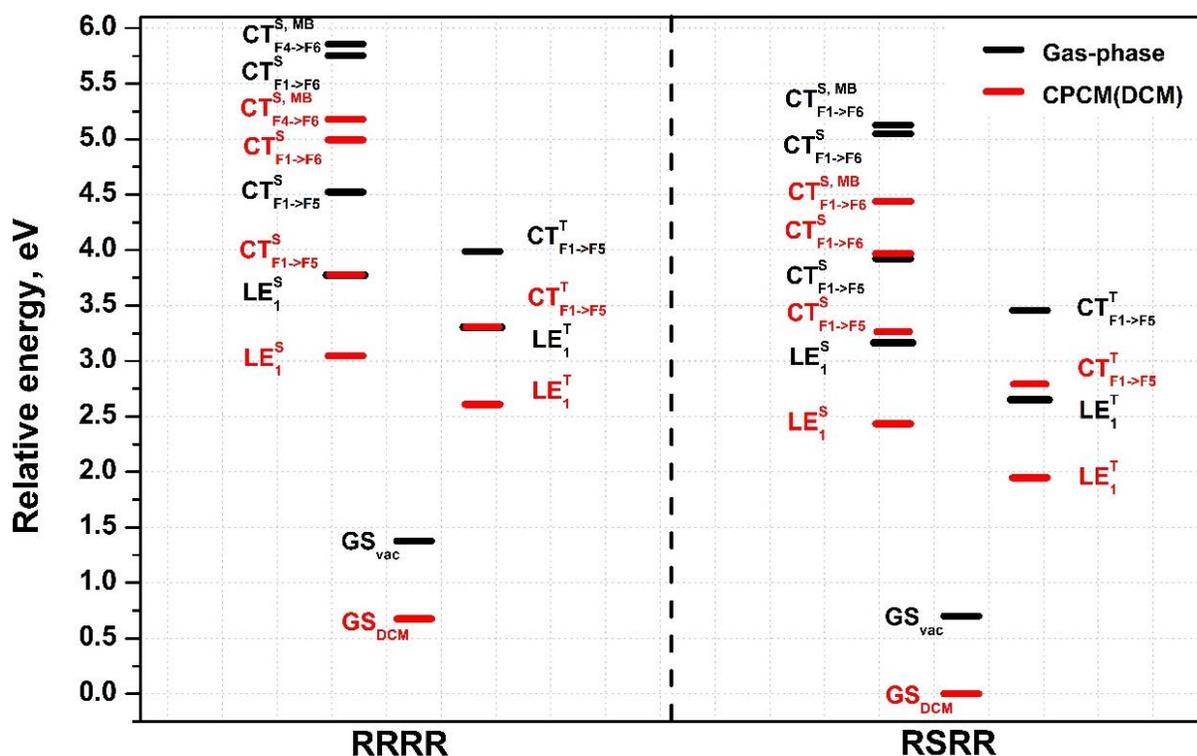


Figure 4. Locally and charge separated excited states energies of Ir pyrrolidino[3,4:1,2][60]fullerene half-sandwich RRRR and RSRR complexes. GS – ground state, LE_1 – lowest locally excited state, $CT_{F1 \rightarrow F5}$ – charge separated excited states corresponded to $F1 \rightarrow F5$ transitions, $CT_{F1 \rightarrow F6}$ – charge separated excited states corresponded to $F1 \rightarrow F6$ transitions, and $CT^{MB}_{F_x \rightarrow F_y}$ – charge separated excited states with biggest oscillator strength, corresponded to denoted in subscript transitions. Superscript index refers to singlet or triplet excited state.

Conclusions

Photoinduced charge separated properties of eight stereoisomers of Ir-based half-sandwich [60]fullerene complexes, have been studied in detail using the TDA DFT approach. The effect of solvent on LE and CT states of interest has been also studied. Solvent effect is found not to be crucial for efficient population of singlet and triplet CT states. Unlike other isomers, the RSRR under photoexcitation exhibits an intense transition corresponded to electron transfer from Ir to fullerene moiety, increasing the partial charge on the metal center. For the first time demonstrated stereospecificity towards photoinduced electron transfer paves the way to photo-chemical applications of the transition-metal complexes with chiral centers in stereoselective photocatalysis. We also believe that will stimulate to take a new look on the already known chiral complexes and encourage further experimental and theoretical research of photoinduced properties of such complexes.

Methods

General. The geometry optimization for all stereoisomers was performed employing the DFT BLYP^{28,29} exchange–correlation functional using the resolution of identity approximation (RI, alternatively termed density fitting)³⁰ implemented in the TURBOMOLE 7.0 program.³¹ Vertical excitation energies were calculated using TDA formalism³² with the range-separated functional from Head-Gordon and co-workers

wB97XD³³ using Gaussian 09 (rev. E01).³⁴ Ahlrichs' Def2-SVP basis set.^{35,36} The empirical dispersion D3 correction with Becke–Johnson damping,^{37,38} was employed. All resulted structures visualized with Chemcraft 1.8.³⁹

Analysis of excited states. The quantitative analysis of exciton delocalization and charge separation in the donor–acceptor complexes was carried out using a tool suggested recently by Plasser et al.^{40,41} A key quantity is the parameter Ω :

$$\begin{aligned}\Omega(A, B) &= \frac{1}{2} \sum_{\alpha \in A, \beta \in B} \left[(SP^{0i})_{\alpha\beta} (P^{0i}S)_{\alpha\beta} + P_{\alpha\beta}^{0i} (SP^{0i}S)_{\alpha\beta} \right] \\ X(F_i) &= \sum_{A \in F_i} \Omega(A, A) \\ CT^{F_i \rightarrow F_j} &= \sum_{A \in F_i, B \in F_j} \Omega(A, B) + \Omega(B, A) \\ CS^{F_i \rightarrow F_j} &= \sum_{A \in F_i, B \in F_j} \Omega(A, B) - \Omega(B, A)\end{aligned}\tag{1}$$

where P^{0i} is the transition density matrix for the $\psi_0 \rightarrow \psi_i$ excitation and S is the overlap matrix. $X(F_i)$ is the extent of exciton localization on the site F_i . $CT^{F_i \rightarrow F_j}$ is the weight of CT configuration in the excited state ψ_i . CT means the total amount of the electron density transferred between fragments, while the CS is the amount of the electron density transferred from one fragment to another. Note that in the situation when charge transfer ($F_i \rightarrow F_j$) is equal to the back transfer ($F_j \rightarrow F_i$) there is no charge separation between the fragments, $CS^{F_i \rightarrow F_j}$ is equal to zero.

Solvent Effects. The equilibrium solvation energy E_s^{eq} in a medium with dielectric constant ϵ was estimated using a COSMO-like polarizable continuum model (C-PCM) in the monopole approximation.⁴²

$$-E_s^{eq}(\mathbf{Q}, \epsilon) = -\frac{1}{2} f(\epsilon) \mathbf{Q}^+ \mathbf{D} \mathbf{Q}, \tag{2}$$

where the $f(\epsilon)$ is the dielectric scaling factor, $f(\epsilon) = \frac{\epsilon - 1}{\epsilon}$, \mathbf{Q} -the vector of n atomic charges in the molecular system, \mathbf{D} is the $n \times n$ symmetric matrix determined by the shape of the boundary surface between solute and solvent; $\mathbf{D} = \mathbf{B}^+ \mathbf{A}^{-1} \mathbf{B}$, where the $m \times m$ matrix \mathbf{A} describes electrostatic interaction between m surface charges and the $m \times n$ \mathbf{B} matrix describes the interaction of the surface charges with n atomic charges of the solute. Atomic charges in the excited state ψ_i , were calculated using Eq. (1).

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